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# Physics in extra dimensions: lecture #1

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*Lecture 1:* **Field theory in compact dimensions.**

**Bosons in the bulk and their collider signatures.**

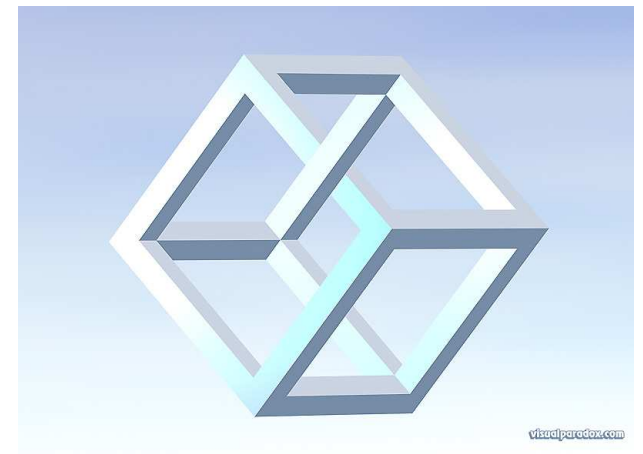
*Lecture 2:* **One universal extra dimension.**

**Discrete symmetries and cascade decays at colliders.**

*Lecture 3:* **6D Standard Model.**

*Lecture 4:* **Particles in a warped extra dimension.**

**Strong coupling at the TeV scale.**



Energy

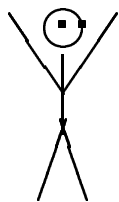


??

~ 1 TeV ?

New Physics

~ 100 GeV



Standard Model

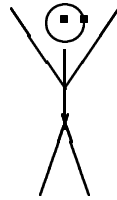
Energy



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New Physics



~ 100 GeV

*Gauge and flavor sectors of the  
Standard Model*

*very weakly interacting particles???*

**“New physics” at the TeV scale could change the basic hypotheses of the Standard Model:**

**local quantum field theory  
in 3 spatial + 1 time dimensions,  
invariant under  $SO(3,1)$  Lorentz transformations.**

*... “terra incognita” ... “uncharted waters” ...*

## Evidence that we live in 3 spatial dimensions:

- it is obvious! (*end of story?!?*)
- Gauss law, in  $3 + n$  spatial dimensions:  $V(r) \sim 1/r^{n+1}$   
We observe  $n = 0$  for gravity and electromagnetism.
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## Counter-arguments:

- what's obvious may be due to preconception  
(*e.g.*, quantum mechanics is not obvious)
- Gauss law may change at short distance
- Standard Model has not been tested below  $10^{-16}$  cm
- gravitational interactions are non-renormalizable in  $D = 3 + 1$

## Bosons in compact spatial dimensions

4D flat spacetime  $\perp$  one dimension of size  $L = \pi R$ :



A scalar field in the bulk,  $\phi(x^\alpha)$ ,  $\alpha = 0, 1, \dots, 4$ :

$$\mathcal{L}_{5D} = (\partial^\mu \phi)^\dagger \partial_\mu \phi - (\partial^4 \phi)^\dagger \partial_4 \phi - m_0^2 \phi^\dagger \phi, \quad \mu = 0, 1, 2, 3$$

$\Rightarrow$  Equation of motion:  $(\partial^\mu \partial_\mu - \partial^4 \partial_4) \phi = m_0^2 \phi$

$m_0$  is the 5D mass of  $\phi$ .

Neumann boundary conditions for “even” fields:

$$\frac{\partial}{\partial x^4}\phi(x^\mu, 0) = \frac{\partial}{\partial x^4}\phi(x^\mu, \pi R) = 0$$



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$$\phi(x^\mu, x^4) = \frac{1}{\sqrt{\pi R}} \left[ \phi^{(0)}(x^\mu) + \sqrt{2} \sum_{j \geq 1} \phi^{(j)}(x^\mu) \cos \left( \frac{jx^4}{R} \right) \right]$$

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Kaluza-Klein  
decomposition

Zero-mode  
(wave function is  
constant along  $x^4$ )

Kaluza-Klein modes:  
particles of definite  
momentum along  $x^4$

4D point of view: a tower of massive particles:

$$\mathcal{L}_{4D} = \int_0^{\pi R} dx^4 \mathcal{L}_{5D} \quad \Rightarrow \quad m_j^2 = m_0^2 + \frac{j^2}{R^2}$$



Dirichlet boundary conditions for “odd” fields:

$$\phi(x, 0) = \phi(x, \pi R) = 0$$

KK decomposition:

$$\phi(x^\mu, x^4) = \frac{\sqrt{2}}{\sqrt{\pi R}} \sum_{j \geq 1} \phi^{(j)}(x^\mu) \sin\left(\frac{jx^4}{R}\right)$$

There is no zero-mode.

The lightest KK mode is  $\phi^{(1)}$ , of mass  $\sqrt{1/R^2 + m_0^2}$

*Homework: Check that the normalization condition for KK functions requires the factor of  $\sqrt{2}$ .*

*Why  $j < 0$  is not allowed?*

## Gauge bosons in 5D:

$A_\mu(x^\nu, x^4)$ ,  $\mu, \nu = 0, 1, 2, 3$ , and

$A_4(x^\nu, x^4)$  – polarization along the extra dimension.

From the point of view of the 4D theory:

$A_4(x^\nu, x^4)$  is a tower of spinless KK modes.

Gauge invariance requires  $A_\mu$  to have a zero-mode:

$$\partial_4 A_\mu(x^\nu, 0) = \partial_4 A_\mu(x^\nu, \pi R) = 0$$

$$A_\mu(x^\nu, x^4) = \frac{1}{\sqrt{\pi R}} \left[ A_\mu^{(0)}(x^\nu) + \sqrt{2} \sum_{j \geq 1} A_\mu^{(j)}(x^\nu) \cos\left(\frac{jx^4}{R}\right) \right]$$

Dirichlet B.C :  $A_4(x^\nu, 0) = A_4(x^\nu, \pi R) = 0$

KK decomposition :  $A_4(x^\nu, x^4) = \sqrt{\frac{2}{L}} \sum_{j \geq 1} A_G^{(j)}(x) \sin\left(\frac{jx^4}{R}\right)$

→  $A_4(x^\nu, x^4)$  does not have a 0-mode! (Odd field)

## Kaluza-Klein spectrum of gauge bosons

$A_G^{(j)}(x^\nu)$  becomes the longitudinal degree of freedom of the spin-1 KK mode  $A_\mu^{(j)}(x^\nu)$ .

$$\begin{array}{ccc} & \vdots & \vdots \\ A_\mu^{(3)} & \text{---} \frac{3}{R} \text{---} & A_G^{(3)} \\ A_\mu^{(2)} & \text{---} \frac{2}{R} \text{---} & A_G^{(2)} \\ A_\mu^{(1)} & \text{---} \frac{1}{R} \text{---} & A_G^{(1)} \\ A_\mu^{(0)} & \text{---} & \end{array}$$

Extra dimensions may be classified according to:

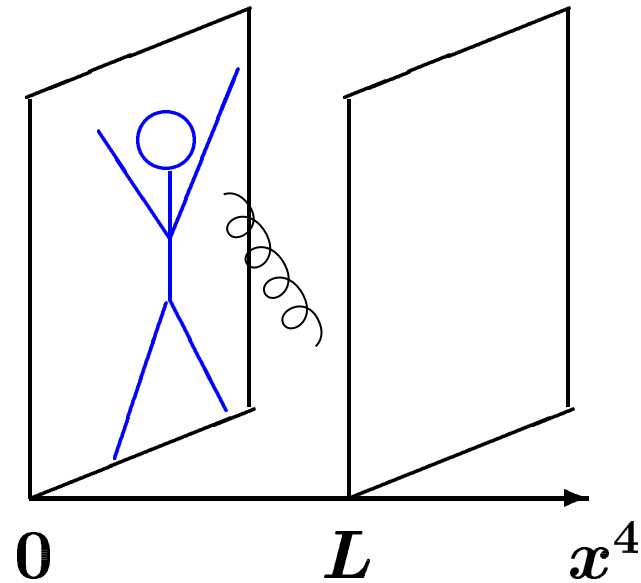
- number (1, 2, ... , 13?)
- type of compactification (*i.e.* boundary conditions)
- metric (flat, warped, ...)
- which fields propagate in the bulk (graviton, top quark, ...)
- existence of localized operators, stabilization mechanism, ...

## Types of extra dimensions:

- graviton only propagates in  $n \geq 2$  flat extra dimensions (ADD)
- bosons only propagate in some flat extra dimensions (DDG)
- bosons and some fermions propagate in flat extra dimensions
- all particles propagate in some flat extra dimensions (UED)
- graviton only propagates in a warped extra dimension (RS)
- all particles propagate in a warped extra dimension
- ...



Assume that only bosons propagate in a flat extra dimension, and that the fermions are localized at  $x^4 = 0$ .

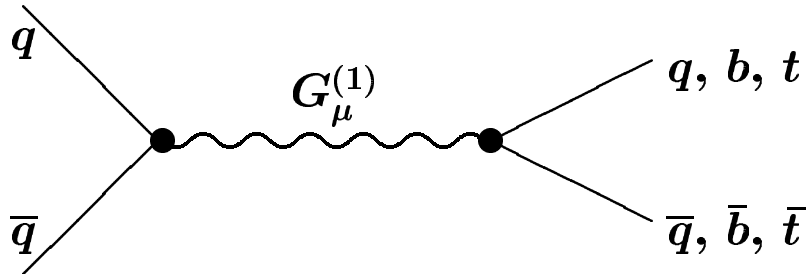


Interactions of the KK gluons with quarks:

$$\begin{aligned} \mathcal{L}_{4D} &= \int_0^L dx^4 g_5 G_\mu^a(x^\nu, x^4) \left[ \delta(x^4) \bar{q}(x^\nu) \gamma^\mu T^a q(x^\nu) \right] \\ &= g_s \left( G_\mu^{(0)a} + \sqrt{2} \sum_{j \geq 1} G_\mu^{(j)a}(x) \right) \bar{q} \gamma^\mu T^a q \end{aligned}$$

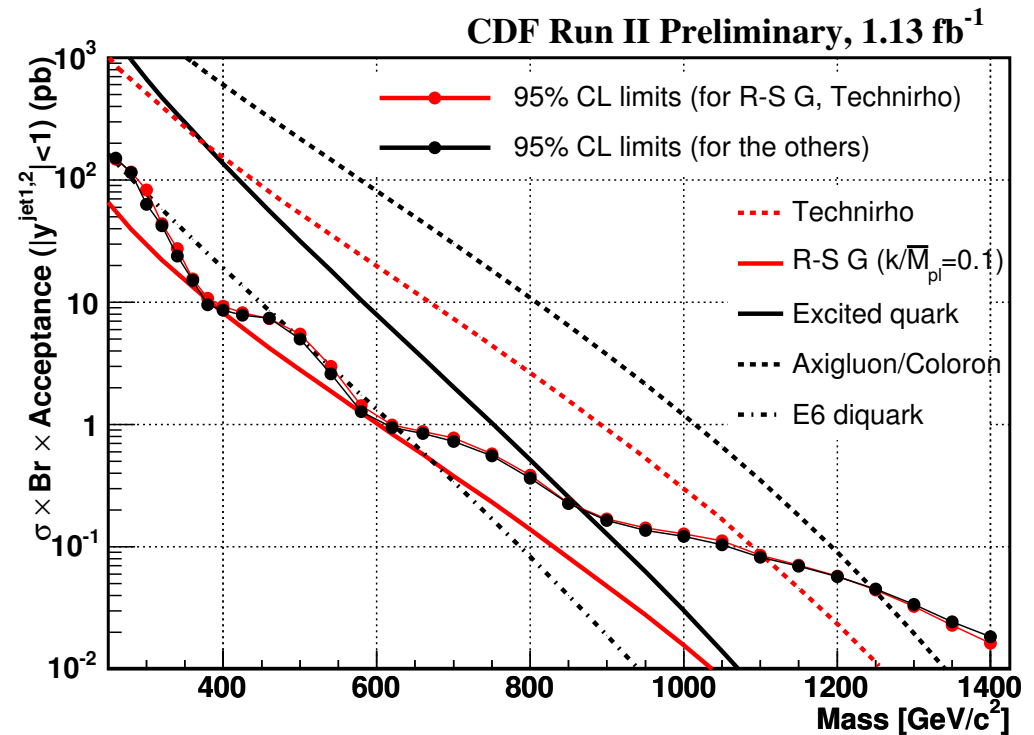
# KK gluon production (in the narrow width approximation):

$$\sigma\left(p\bar{p} \rightarrow G_{\mu}^{(1)} X\right) \approx \frac{16\pi^2\alpha_s}{9s} \sum_q \int_{M^2/s}^1 \frac{dx}{x} \left[ q(x) q\left(\frac{M^2}{xs}\right) + \bar{q}(x) \bar{q}\left(\frac{M^2}{xs}\right) \right]$$



Run II limit on dijet resonances:

$$\longrightarrow 1/R > 1.4 \text{ TeV}$$



*5D theory = 4D theory with some heavy particles*

**$SU(3)_c$  in extra dimensions  $\rightarrow$  SM gluon + heavy gluons**

**4D theory with the same spectrum must include a larger gauge symmetry.**

**$SU(3)_1 \times SU(3)_2 \rightarrow SU(3)_c$  spontaneously broken by the VEV  
of a scalar transforming as  $(3, \bar{3})$**

**Quarks transform as 3 of  $SU(3)_1$**

$G_\mu^a$  - massless gluon as in QCD, with  $g_s = \frac{h_1 h_2}{\sqrt{h_1^2 + h_2^2}}$

$G_\mu'^a$  - massive gluon (“coloron”) with couplings  $g_s \frac{h_1}{h_2} G_\mu'^a \bar{q} \gamma^\mu T^a q$

KK gluon coupling recovered if the two  $SU(3)$  gauge couplings satisfy  $h_1/h_2 = \sqrt{2}$ .

The 4D theory describing the first  $N$  KK modes of the gluon has a  $SU(3)_1 \times SU(3)_2 \times \cdots \times SU(3)_{N+1} \rightarrow SU(3)_c$  gauge structure.

**Assume that the Higgs doublet and electroweak gauge bosons propagate in one flat extra dimension, while the fermions are localized at one end of the interval.**

*Homework: Derive the couplings of the KK electroweak gauge bosons to quark and leptons.*

**LEP-II limits on four-fermion interactions imply  $1/R > 6$  TeV.**

## Fermions in 5D:

Chiral boundary conditions:

$$\begin{aligned}\chi_L(x^\mu, 0) &= \chi_L(x^\mu, \pi R) = 0 \\ \frac{\partial}{\partial x^4} \chi_R(x^\mu, 0) &= \frac{\partial}{\partial x^4} \chi_R(x^\mu, \pi R) = 0\end{aligned}$$

Kaluza-Klein decomposition:

$$\chi = \frac{1}{\sqrt{\pi R}} \left\{ \chi_R^0(x^\mu) + \sqrt{2} \sum_{j \geq 1} \left[ \chi_R^j(x^\mu) \cos \left( \frac{\pi j x^4}{L} \right) + \chi_L^j(x^\mu) \sin \left( \frac{\pi j x^4}{L} \right) \right] \right\}$$

*Something to read:*

**G. D. Kribs, TASI lectures on the  
“Phenomenology of extra dimensions”, hep-ph/0605325.**

**E. H. Simmons, “Coloron phenomenology,”  
Phys. Rev. D 55, 1678 (1997), hep-ph/9608269.**

**K. R. Dienes, E. Dudas and T. Gherghetta, “Grand unification  
at intermediate mass scales through extra dimensions,”  
Nucl. Phys. B 537, 47 (1999), hep-ph/9806292.**

...

## Conclusions so far

- Extra spatial dimensions may exist if their size is small enough. The limits are model dependent.
- Any particle that propagates in  $D \geq 5$  would appear in experiments as a tower of heavy 4-dimensional particles.
- If bosons propagate in extra dimensions while fermions are localized, then the Kaluza-Klein bosons may be singly produced, leading to  $s$ -channel resonances.
- Kaluza-Klein modes of the quarks and leptons are vectorlike fermions.